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Some properties of yarn

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Dr J. Th. Runnenburg

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### 1. Introduction

A yarn, of nonnegative cross-section, lies along a t-axis with  $-\infty < t < \infty$ . It is composed of fibres, which are cylinders with their axes parallel to the t-axis. The head of a fibre is its left endpoint on the t-axis, the tail its right endpoint. Tibre heads are distributed along the t-axis according to a Poisson-process with parameter  $\lambda(t)$ , where  $\lambda(t)$  is Lebesgue-integrable on any finite interval and  $0 < \lambda(t) \le \lambda$  for all t and some finite  $\lambda$ . We define such a process in the following way: Regarding the t-axis as composed of intervals, of which the j<sup>th</sup> interval is (j,j+1] (open on the left, closed on the right), where j runs through the integers, the distribution of the number  $r_j$  of heads falling in the j<sup>th</sup> interval is given by

(1.1) 
$$P\{\underline{r}_{j} = r\} = e^{-\lambda} j - \frac{\lambda j}{r!} \quad \text{for } r \ge 0$$

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(1.2) 
$$\lambda_{j} \stackrel{\text{def}}{=} \int_{j}^{j+1} \lambda(t) dt.$$

The head of the  $k^{th}$  fibre  $(k=1,2,\ldots,r;\ \underline{r_j}=r\geq 1)$  having its head in the  $j^{th}$  interval lies in  $\underline{t_j},k$ , where under the condition  $\underline{r_j}=r$  the  $\underline{t_j},1,\ldots,\underline{t_j},r$  are independently distributed random variables with the common distribution function  $K_j(t)$ , with

$$(1.3) F_{j}(t) \stackrel{\text{def}}{=} F\{\underline{t}_{j,1} \le t\} = \begin{cases} 0 & \text{for } t \le j, \\ \lambda_{j}^{-1} \int_{j}^{t} \lambda(u) du & \text{for } j \le t \le j+1, \\ 1 & \text{for } j+1 \le t. \end{cases}$$

The  $\underline{t}_{j,k}$  for different intervals on the t-axis are mutually independent. The length  $\underline{x}$  ( $\geq$  0) and cross-section  $\underline{y}$  ( $\geq$  0) of a fibre have a simultaneous distribution function  $H(x,y) = \mathbb{P}\{\underline{x} \leq x, \ \underline{y} \leq y\}$ . The vectors ( $\underline{x},\underline{y}$ ) belonging to different fibres are mutually independent and independent of the location of the fibre on the t-axis. A realization of the yarn is a stepfunction, which can be taken continuous from the right. The length of the kth fibre having its head in the jth interval is  $\underline{x}_{j,k}$ , its cross-section  $\underline{y}_{j,k}$ .

Object of this paper is the study of random variables like  $\underline{c}(t_0)$  and  $\underline{v}(t_0,t_0+h)$ , which denote the cross-section of the yarn at  $t_0$  and the volume of yarn in the interval  $(t_0,t_0+h]$  respectively (with  $t_0$  arbitrary and h>0).

## 2. Terpstra's approach

In this section we assume that  $\mathbb{H}(x_0,\infty)=1$  for some  $x_0<\infty$ , i.e. the fibres have length  $\leq x_0$  with probability 1. Furthermore we take  $\lambda(t)=\lambda>0$ , i.e. we consider a stationary Poisson-process (this restriction need not be made).

In order to derive  $\mathcal{E}_{e}^{i\tau c(t_{0})}$ , we first compute

$$(2.1) \qquad \varphi_{\underline{c}}(\tau, t_0; s, 1) \stackrel{\text{def}}{=} \{ e^{i\tau \underline{c}(t_0)} | \underline{n}(s, s+1) = 1 \},$$

where  $\underline{n}(s,s+1)$  denotes the number of heads of fibres in the interval (s,s+1] and  $s < t_0 - x_0 < t_0 < s+1$ .

If  $\underline{t}$  is the coordinate of the head of the one fibre falling in (s,s+1], then it is known that  $\underline{t}$  is uniformly distributed in this interval. If the fibre has length  $\underline{x}=x$ , with probability  $\frac{x}{l}$  the head lies in  $(t_0-x,t_0]$ . Hence

$$(2.2) \qquad \varphi_{\underline{c}}(\tau, t_0; s, 1) = \int_0^\infty \int_0^\infty \left[ e^{i\tau y} \frac{x}{1} + e^0 \left( 1 - \frac{x}{1} \right) \right] dH(x, y) =$$

$$= 1 - \frac{1}{1} \int_0^\infty \int_0^\infty x(1 - e^{i\tau y}) dH(x, y).$$

As

(2.3) 
$$P\{\underline{n}(s, s+1) = n\} = e^{-\lambda 1} \frac{(\lambda_1)^n}{n!}$$
 for  $n = 0, 1, 2, ...$ 

and  $\underline{c}(t_0)$  is the sum of the independent contributions to the cross-section of all fibres with heads in (s, s+1], we have (the characteristic function of a sum of independent random variables being equal to the product of the characteristic functions of the individual terms)

(2.4) 
$$\varphi_{\underline{c}}(\tau, t_{o}) \stackrel{\text{def}}{=} \begin{cases} e^{i\tau} \underline{c}(t_{o}) \\ e^{-\lambda t_{o}} \end{cases} =$$

$$= \sum_{n=0}^{\infty} \{ \psi_{\underline{c}}(\tau, t_{o}; s, 1) \}^{n} e^{-\lambda t_{o}} \frac{(\lambda t_{o})^{n}}{n!} =$$

$$= \exp(-\lambda t_{o}) \left\{ 1 - (1 - \frac{1}{t_{o}}) \int_{0}^{\infty} x(1 - e^{i\tau y}) dH(x, y) \right\} =$$

$$= \exp(-\lambda t_{o}) \int_{0}^{\infty} x(1 - e^{i\tau y}) dH(x, y) .$$

As was to be expected,  $\varphi_{\underline{C}}(\tau, t_0)$  does not depend on either s or 1. We may conjecture, that (2.4) will also hold if  $H(x, \infty) < 1$  for all  $x < \infty$ . (Some restriction is needed: we must have  $\underbrace{\mathcal{E}}_{\underline{X}} < \infty$ , cf. Breny (1957) and section 4.)

We remark that from (2.4) it is evident that  $\underline{c}(t_0)$  has a compound Poisson - distribution, i.e.

(2.5) 
$$\underline{c}(t_0) = \frac{\underline{n}}{\underline{b}} \underline{b}_{\underline{j}},$$

where

(2.6) 
$$P\{\underline{n} = n\} = e^{-\lambda \frac{\mathcal{E}}{\underline{x}}} \frac{(\lambda \frac{\mathcal{E}}{\underline{x}})^n}{n!} \quad \text{for } n = 0, 1, 2, \dots$$

and, under the condition  $\underline{n} = n$ , the  $\underline{b}_j$  are independently distributed, each with distribution function

(2.7) 
$$G(y) \stackrel{\text{def}}{=} o^{\int_{-\infty}^{\infty} x \, d_{\mathbf{x}} H(x, y)} \underbrace{\mathcal{E}_{\underline{\mathbf{x}}}}.$$

Analogously to the derivation of (2.4) one can show that

(2.8) 
$$\varphi_{\underline{v}}(\tau, t_{o}, h) \stackrel{\text{def}}{=} \mathcal{E}_{e}^{i\tau \underline{v}(t_{o}, t_{o} + h)} =$$

$$= \exp -\lambda \left[h + \mathcal{E}_{\underline{x}} + \int_{y=0}^{\infty} \int_{x=0}^{h} \left\{ \frac{2}{iy\tau} \left(1 - e^{i\tau xy}\right) - (h - x) e^{i\tau xy} \right\} dH(x, y) +$$

$$+ \int_{y=0}^{\infty} \int_{x=h}^{\infty} \frac{2}{iy\tau} \left(1 - e^{i\tau hy}\right) - (x - h) e^{i\tau hy} dH(x, y) \right],$$

always assuming a finite maximum length  $x_0$  for the fibres.

Both (2.4) and (2.8) are due to Terpstra. These results are known for  $\underline{y} = 1$  with probability 1 (cf. Spencer - Smith and Todd (1941), Martindale (1945), Breny (1952) and Breny (1953), Olerup (1952)). Related considerations as to method of derivation are to be found in Fortet (1951).

#### 3. Van Dantzig's method

In this section we assume that  $\underline{y}\equiv 1$  (or  $\underline{y}=1$  with probability 1). In Breny (1957) relation (2.8) (with  $\underline{y}\equiv 1$ ) is obtained by applying a limiting procedure to (in our notation)  $\mathcal{E}$  exp i  $\sum_{j=0}^{n} \tau_{j} \underline{n}(t_{j})$ , expressed as a complicated sum of double integrals. As Van Dantzig pointed out to Breny (cf. Breny (1957), page 33) one might proceed in the following manner. Let A denote a Lebesgue-measurable set in the space  $\Omega = \{(t, \mathbf{x}) \mid -\infty < t < \infty, 0 \le \mathbf{x} < \infty\}$  and  $\underline{m}(A)$  the number of fibres for which  $(\underline{t}_{j}, \underline{k}, \underline{x}_{j}, \underline{k}) \in A$  is satisfied. Then, if  $A \cap B = 0$ ,  $\underline{m}(A)$  and  $\underline{m}(B)$ 

are independent stochastic variables having a Poisson - distribution, the parameter of m(A) being

(3.1) 
$$\{ \underline{m}(A) = \iint_{A} \lambda(t) dt dF(x) ,$$

where  $F(x) \stackrel{\text{def}}{=} P\{\underline{x} \le x\}$  is the distribution function of the length of a gibre. For a general class of real functions  $\xi(t,x)$  one can define

(3.2) 
$$\int_{0}^{\infty} \int_{-\infty}^{\infty} \xi(t,x) d\underline{m},$$

where <u>m</u> is a <u>stochastic measure</u> on  $\Omega$ , with  $\underline{m}(A)$  as described, for every L-measurable set A. Because of the independence of the  $\underline{m}(A)$  for disjoint sets, we shall have something like (use Riemann-sums)

$$(3.3) \qquad \text{E exp it } \int_{0}^{\infty} \int_{-\infty}^{\infty} \xi(t,x) \, d\underline{m} \approx \prod_{\nu=0}^{\infty} \prod_{\mu=-\infty}^{\infty} \sum_{k=0}^{\infty} \frac{i^{\tau} \xi(t_{\mu}^{*}, x_{\nu}^{*}) \underline{m}(\Delta_{\mu\nu})}{k^{\tau} \sum_{\nu=0}^{\infty} \prod_{\mu=-\infty}^{\infty} \left\{ (1 - \lambda(t) \, \Delta t \, \Delta F(x)) e^{0} + \lambda(t) \, \Delta t \, \Delta F(x) e^{i\tau \xi(t_{\mu}^{*}, x_{\nu}^{*})} \right\} \approx \sum_{\nu=0}^{\infty} \prod_{\mu=-\infty}^{\infty} \exp - \lambda(t) \{ 1 - i^{\tau} \xi(t_{\mu}^{*}, x_{\nu}^{*}) \} \Delta t \, \Delta F(x) \approx \sum_{\nu=0}^{\infty} \prod_{\mu=-\infty}^{\infty} \sum_{n=0}^{\infty} \sum_{n=-\infty}^{\infty} \lambda(t) \{ 1 - i_{\tau} \xi(t, x) \} \, dt \, dF(x)$$

or

(3.4) 
$$\mathcal{E} \exp i\tau \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} \xi(t,x) \, d\underline{m} = \exp - \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} \lambda(t) \{1 - i\tau \xi(t,x)\} \, dt \, dF(x) .$$

Substitution of

(3.5) 
$$\xi(t,x) = \begin{cases} 1 & \text{for } t \leq t_0, t+x > t_0 \\ 0 & \text{otherwise} \end{cases},$$

now leads to (2.4), because

$$\underline{\mathbf{c}}(\mathbf{t}_{0}) = \int_{0}^{\infty} \int_{\mathbf{t}_{0}-\mathbf{x}<\mathbf{t}\leq\mathbf{t}_{0}} d\underline{\mathbf{m}}.$$

In the same way

$$\xi(t,x) = \{(t-t_0) \cdot (t_0-t) - (t-t_0-h) \cdot (t_0+h-t)\} + \{(x+t-t_0) \cdot (t_0-x-t) - (x+t-t_0-h) \cdot (t_0+h-x-t)\}$$

leads to (2.8).

Stochastic set functions have been discussed by Prékopa (cf. Prékopa (1956, 1957)).

# 4. Main formula

The  $k^{th}$  fibre from the j<sup>th</sup> interval contributes to the cross-section at **point** t of the axis an amount

$$(4.1) \qquad \underline{c}_{j,k}(t) \stackrel{\text{def}}{=} y_{j,k} \{ \iota(\underline{t}_{j,k} + \underline{x}_{j,k} - t) - \iota(\underline{t}_{j,k} - t) \},$$

where l(x) = 1 for  $x \ge 0$  and l(x) = 0 for x < 0. Hence we find for the total cross - section at point  $t^{-1}$ 

(4.2) 
$$\underline{\underline{c}}(t) \stackrel{\text{def}}{=} \sum_{j=-\infty}^{\infty} \sum_{k=1}^{\underline{r}j} \underline{y}_{j,k} \{ i(\underline{t}_{j,k} + \underline{x}_{j,k} - t) - i(\underline{t}_{j,k} - t) \}.$$

We shall study the random variable

(4.3) 
$$\underline{u}_{T} \stackrel{\text{def}}{=} \underbrace{\overset{\infty}{c}}(t) dT(t)$$

under the condition  $\mathcal{E}_{\underline{x}} < \infty$ . This variable is (as will be shown) well-defined for any real-valued function T(t), which is of bounded variation in a closed interval  $\begin{bmatrix} t_1, t_2 \end{bmatrix}$  (with  $-\infty < t_1 < t_2 < \infty$ ), constant for  $t \leq t_1$ , as well as constant for  $t \geq t_2$  and finally continuous from the right for all t. We assume (this is no restriction) that  $t_1$  is a negative integer and  $t_2$  a positive integer.

With probability 1  $\underline{c}(t)$  is a stepfunction with a finite number of finite steps in the interval  $[t_1,t_2]$ , given  $\{\underline{x} < \infty \text{ (Breny (1957))}.$  As inside  $[t_1,t_2]$  only a finite number of finite steps originate, it is sufficient to prove that with probability 1 only a finite number of fibres cover the point  $t_1$ . Thus we have to prove

or

We take  $\sum_{k=1}^{\underline{r}_j} \dots = 0$  for  $\underline{r}_j = 0$ , here and later.

$$(4.5) P\left\{\max_{1\leq k\leq r_j} \left(t_{j,k} + x_{j,k} - t_1\right) \geq 0 \text{ for infinitely many } j\right\} = 0.$$

To prove the last relation, it is sufficient to show (by one of the Borel - Cantelli lemmas) that for the (independent) events

(4.6) 
$$\underset{j}{\text{A}} = \left\{ \begin{array}{l} \max & (\underline{t}_{j,k} + \underline{x}_{j,k} - \underline{t}_{1}) \ge 0 \\ 1 \le k \le \underline{r}_{j} \end{array} \right.$$

we have

$$\begin{array}{c}
t_1 - 1 \\
\Sigma \\
j = -\infty
\end{array} P\{A_j\} < \infty.$$

Now, if  $F(x) \stackrel{\text{def}}{=} P\{\underline{x} \leq x\}$ ,

as  $1-e^{-x} \le x$  for  $x \ge 0$  and  $\lambda_j \le \lambda$  for each j. Hence, because  $\sum\limits_{k=1}^{\infty} \left\{1-F(k)\right\} \le \sum\limits_{0}^{\infty} x \, dF(x) = \underbrace{ \sum\limits_{x} < \infty}_{0}$ , condition (4.7) is satisfied.

For a stochastic stepfunction  $\underline{c}(t)$  which has, with probability 1, a finite number of finite steps in the interval  $\begin{bmatrix} t_1, t_2 \end{bmatrix}$ , the integral in (4.3) is defined with probability 1 as the Lebesgue-Stieltjes integral of the realization c(t) and satisfies

$$\underline{u}_{T} = \int_{-\infty}^{\infty} \underline{c}(t) dT(t) =$$

$$= \int_{j=-\infty}^{\infty} \frac{\underline{r}_{j}}{k=1} \underline{y}_{j,k} \int_{-\infty}^{\infty} \{ \underline{t}(\underline{t}_{j,k} + \underline{x}_{j,k} - t) - \underline{i}(\underline{t}_{j,k} - t) \} dT(t) =$$

$$= \int_{j=-\infty}^{\infty} \frac{\underline{r}_{j}}{k=1} \underline{y}_{j,k} \{ \underline{T}(\underline{t}_{j,k} + \underline{x}_{j,k}) - \underline{T}(\underline{t}_{j,k}) \} ,$$

where the double series has only a finite number of nontrivial terms (i.e. terms unequal 0) with probability 1.

Lemma. If  $\underline{a}_n$ ,  $n=1,2,\ldots$  are mutually independent real-valued stochastic variables, such that the events  $A_n = \{\underline{a}_n \neq 0\}$  satisfy  $P\{A_n \text{ occurs for infinitely many } n\} = 0$ , we have:  $\underline{s} = \sum_{n=1}^{\infty} \underline{a}_n$  is a stochastic variable for which (for all real  $\tau$ )

$$(4.10) \qquad \qquad \qquad \mathcal{E}_{e}^{i\tau \underline{s}} = \prod_{n=1}^{\infty} \mathcal{E}_{e}^{i\tau \underline{a}\underline{n}}.$$

Proof; s is a well-defined stochastic variable, for it is with probability 1 the sum of a finite number of independent random variables unequal 0.

As we have

(4.11) 
$$\{A_n \text{ occurs for infinitely many } n\} = \bigcap_{n=1}^{\infty} \bigcup_{m=n}^{\infty} A_m$$

and therefore

(4.12) 
$$\lim_{n\to\infty} P\{\bigcup_{m=n}^{\infty} A_m\} = 0,$$

to each  $\epsilon > 0$  there exists an  $N = N(\epsilon)$ , for which  $P\{\bigcup_{m=N+1}^{\infty} A_m\} < \epsilon$  or  $P\{\bigcap_{m=N+1}^{\infty} \overline{A}_m\} > 1 - \epsilon$  (if  $\overline{A}_m$  denotes the complement of the set  $A_m$ ). But then

$$(4.13) \qquad \left| \underbrace{\sum_{n=1}^{\infty} e^{i\tau \underline{a}_{n}} - \underbrace{\sum_{n=1}^{N} e^{i\tau \underline{a}_{n}}}_{i\tau \underline{a}_{n}} \right| \leq \left| \underbrace{\sum_{n=1}^{N} e^{i\tau \underline{a}_{n}}}_{i\tau \underline{a}_{n}} \right| \cdot \left| \underbrace{\sum_{n=N+1}^{\infty} e^{i\tau \underline{a}_{n}}}_{i\tau \underline{a}_{n}} - 1 \right| \leq \left| \underbrace{\sum_{n=N+1}^{\infty} e^{i\tau \underline{a}_{n}}}_{i\tau \underline{a}_{n}} - 1 \right| \leq \left| \underbrace{\sum_{n=N+1}^{\infty} e^{i\tau \underline{a}_{n}}}_{i\tau \underline{a}_{n}} - 1 \right| \leq \left| \underbrace{\sum_{n=N+1}^{\infty} e^{i\tau \underline{a}_{n}}}_{i\tau \underline{a}_{n}} - 1 \right| \leq \left| \underbrace{\sum_{n=N+1}^{\infty} e^{i\tau \underline{a}_{n}}}_{i\tau \underline{a}_{n}} - 1 \right| \leq \left| \underbrace{\sum_{n=N+1}^{\infty} e^{i\tau \underline{a}_{n}}}_{i\tau \underline{a}_{n}} - 1 \right| \leq \left| \underbrace{\sum_{n=N+1}^{\infty} e^{i\tau \underline{a}_{n}}}_{i\tau \underline{a}_{n}} - 1 \right| \leq \left| \underbrace{\sum_{n=N+1}^{\infty} e^{i\tau \underline{a}_{n}}}_{i\tau \underline{a}_{n}} - 1 \right| \leq \left| \underbrace{\sum_{n=N+1}^{\infty} e^{i\tau \underline{a}_{n}}}_{i\tau \underline{a}_{n}} - 1 \right| \leq \left| \underbrace{\sum_{n=N+1}^{\infty} e^{i\tau \underline{a}_{n}}}_{i\tau \underline{a}_{n}} - 1 \right| \leq \left| \underbrace{\sum_{n=N+1}^{\infty} e^{i\tau \underline{a}_{n}}}_{i\tau \underline{a}_{n}} - 1 \right| \leq \left| \underbrace{\sum_{n=N+1}^{\infty} e^{i\tau \underline{a}_{n}}}_{i\tau \underline{a}_{n}} - 1 \right| \leq \left| \underbrace{\sum_{n=N+1}^{\infty} e^{i\tau \underline{a}_{n}}}_{i\tau \underline{a}_{n}} - 1 \right| \leq \left| \underbrace{\sum_{n=N+1}^{\infty} e^{i\tau \underline{a}_{n}}}_{i\tau \underline{a}_{n}} - 1 \right| \leq \left| \underbrace{\sum_{n=N+1}^{\infty} e^{i\tau \underline{a}_{n}}}_{i\tau \underline{a}_{n}} - 1 \right| \leq \left| \underbrace{\sum_{n=N+1}^{\infty} e^{i\tau \underline{a}_{n}}}_{i\tau \underline{a}_{n}} - 1 \right| \leq \left| \underbrace{\sum_{n=N+1}^{\infty} e^{i\tau \underline{a}_{n}}}_{i\tau \underline{a}_{n}} - 1 \right| \leq \left| \underbrace{\sum_{n=N+1}^{\infty} e^{i\tau \underline{a}_{n}}}_{i\tau \underline{a}_{n}} - 1 \right| \leq \left| \underbrace{\sum_{n=N+1}^{\infty} e^{i\tau \underline{a}_{n}}}_{i\tau \underline{a}_{n}} - 1 \right| \leq \left| \underbrace{\sum_{n=N+1}^{\infty} e^{i\tau \underline{a}_{n}}}_{i\tau \underline{a}_{n}} - 1 \right| \leq \left| \underbrace{\sum_{n=N+1}^{\infty} e^{i\tau \underline{a}_{n}}}_{i\tau \underline{a}_{n}} - 1 \right| \leq \left| \underbrace{\sum_{n=N+1}^{\infty} e^{i\tau \underline{a}_{n}}}_{i\tau \underline{a}_{n}} - 1 \right| \leq \left| \underbrace{\sum_{n=N+1}^{\infty} e^{i\tau \underline{a}_{n}}}_{i\tau \underline{a}_{n}} - 1 \right| \leq \left| \underbrace{\sum_{n=N+1}^{\infty} e^{i\tau \underline{a}_{n}}}_{i\tau \underline{a}_{n}} - 1 \right| \leq \left| \underbrace{\sum_{n=N+1}^{\infty} e^{i\tau \underline{a}_{n}}}_{i\tau \underline{a}_{n}} - 1 \right| \leq \left| \underbrace{\sum_{n=N+1}^{\infty} e^{i\tau \underline{a}_{n}}}_{i\tau \underline{a}_{n}} - 1 \right| \leq \left| \underbrace{\sum_{n=N+1}^{\infty} e^{i\tau \underline{a}_{n}}}_{i\tau \underline{a}_{n}} - 1 \right| \leq \left| \underbrace{\sum_{n=N+1}^{\infty} e^{i\tau \underline{a}_{n}}}_{i\tau \underline{a}_{n}} - 1 \right| \leq \left| \underbrace{\sum_{n=N+1}^{\infty} e^{i\tau \underline{a}_{n}}}_{i\tau \underline{a}_{n}} - 1 \right| \leq \left| \underbrace{\sum_{n=N+1}^{\infty} e^{i\tau \underline{a}_{n}}}_{i\tau \underline{a}_{n}} - 1 \right| \leq \left| \underbrace{\sum_{n=N+1}^{\infty} e^{i\tau \underline{a}_{n}}}_{i\tau \underline{a}_{n}} - 1 \right| \leq \left| \underbrace{\sum_{n=N+1}^{\infty} e^{i\tau \underline{a}_{n}}}_{i\tau \underline{a}_{n}} - 1 \right| \leq \left| \underbrace{\sum_{n$$

and hence

(4.14) 
$$\lim_{N\to\infty} \prod_{n=1}^{N} \mathcal{E} e^{i\tau \underline{a}_n} = \mathcal{E} \prod_{n=1}^{\infty} e^{i\tau \underline{a}_n},$$

which proves the lemma.

The random variables

(4.15) 
$$\underline{z}_{j} \stackrel{\text{def}}{=} \frac{\underline{r}_{j}}{\prod_{k=1}^{k}} \exp i \tau \underline{y}_{j \neq k} \{ \underline{r}(\underline{t}_{j,k} + \underline{x}_{j,k}) - \underline{r}(\underline{t}_{j,k}) \}$$

are nontrivial, i.e. not equal to 1 with probability 1, only for  $j \le t_2 - 1$ . Hence we find from (4.9) by applying the lemma to

$$\underline{a}_{n} \stackrel{\text{def}}{=} \underline{z}_{t_{2}-n} ,$$

that for all real t

$$(4.17) \quad \mathcal{E} \quad \exp \quad i\tau \underline{u}_{T} =$$

$$= \mathcal{E} \prod_{j=-\infty}^{\infty} \prod_{k=1}^{x_{j}} \exp \quad i\tau \underline{y}_{j,k} \{ T(\underline{t}_{j,k} + \underline{x}_{j,k}) - T(\underline{t}_{j,k}) \} =$$

$$= \prod_{j=-\infty}^{\infty} \mathcal{E} \prod_{k=1}^{x_{j}} \exp \quad i\tau \underline{y}_{j,k} \{ T(\underline{t}_{j,k} + \underline{x}_{j,k}) - T(\underline{t}_{j,k}) \}.$$

Further we have by (1,1), (1.2) and (1.3), using Fubini's theorem to **prove** the second equality,

$$\begin{array}{ll} (\textbf{4.18}) & \mathcal{E} \prod_{k=1}^{\frac{j-1}{2}} \exp i\tau \underline{y}_{j,k} \{ \mathbb{T}(\underline{t}_{j,k} + \underline{x}_{j,k}) - \mathbb{T}(\underline{t}_{j,k}) \} = \\ & = \mathcal{E} \prod_{r=0}^{\infty} \{ \prod_{k=1}^{r} \exp i\tau \underline{y}_{j,k} [\mathbb{T}(\underline{t}_{j,k} + \underline{x}_{j,k}) - \mathbb{T}(\underline{t}_{j,k})] \Big| \underline{r}_{j} = r \} P \{ \underline{r}_{j} = r \} = \\ & = \prod_{r=0}^{\infty} \{ \int_{0}^{\infty} \int_{0}^{\infty} \int_{0}^{j+1} \exp i\tau \underline{y} [\mathbb{T}(t+x) - \mathbb{T}(t)] dK_{j}(t) dH(x,y) \}^{r} P \{ \underline{r}_{j} = r \} = \\ & = \exp \{ -\lambda_{j} (1 - \int_{0}^{\infty} \int_{0}^{\infty} \int_{0}^{j+1} \exp i\tau \underline{y} [\mathbb{T}(t+x) - \mathbb{T}(t)] dK_{j}(t) dH(x,y) ) \} = \\ & = \exp - \int_{0}^{\infty} \int_{0}^{\infty} \int_{0}^{j+1} \lambda(t) [1 - \exp i\tau \underline{y} \{ \mathbb{T}(t+x) - \mathbb{T}(t) \} ] dt dH(x,y) . \end{aligned}$$

Because  $|1 - \exp i\tau y \{T(t+x) - T(t)\}| \le 2$  and T(t+x) - T(t) = 0 for  $t \le t_1 - x$  and  $t \ge t_2$ , we have

(4.19) 
$$\int_{-\infty}^{\infty} |1 - \exp i\tau y \{T(t + x) - T(t)\}| dt \leq 2\{x + (t_2 - t_1)\}$$

and so, because

$$(4.20) \qquad \int_{0}^{\infty} \int_{-\infty}^{\infty} \left| 1 - \exp i\tau y \left\{ T(t+x) - T(t) \right\} \right| dt dH(x,y) \leq 2 \left\{ \frac{x}{2} + \left( t_2 - t_1 \right) \right\} < \infty,$$

we are allowed to apply Fubini's theorem once again. Combining (4.17) and (4.18) we find our main formula

$$(4.21) \qquad \mathcal{E} \exp i\tau \underline{u}_{T} =$$

$$= \exp - \sum_{j=-\infty}^{\infty} \int_{0}^{\infty} \int_{0}^{\infty} \int_{0}^{\infty} \lambda(t)[1 - \exp i\tau y\{T(t+x) - T(t)\}] dt dH(x,y) =$$

$$= \exp - \int_{0}^{\infty} \int_{0}^{\infty} \int_{\infty}^{\infty} \lambda(t)[1 - \exp i\tau y\{T(t+x) - T(t)\}] dt dH(x,y).$$

The particular cases

(4.22) 
$$T(t) = L(t - t_0),$$

$$T(t) = (t - t_0) l(t - t_0) - (t - t_0 - h) l(t - t_0 - h),$$

show, that both (2.4) and (2.8) are satisfied for  $\{\xi,\underline{x}\in\infty\}$ . In fact a generalization of these results (with  $\lambda$  (t) instead of  $\lambda$  ) has been obtained. Relation (4.21) may be used to obtain interesting formulae for other stochastic variables besides  $\underline{c}(t)$  and  $\underline{v}(t)$ , t + h).

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